

Degree Complexity of a Family of Birational Maps

Eric Bedford*, Kyounghee Kim, Truong Trung Tuyen,
Nina Abarenkova†, and Jean-Marie Maillard

§0. Introduction. Birational mappings on the space of $q \times q$ matrices have been found to arise as natural symmetries in lattice statistical mechanics. One such map gives rise to a family $k_{a,b}$ of birational maps of the plane (see [BMR1,2], [A2]). Dynamical properties of this family have been studied in a number of works ([A1–8], [BD2], [BMR], [BM]). Recall the quantity

$$\delta(k) := \lim_{n \rightarrow \infty} (\deg(k^n))^{\frac{1}{n}},$$

which is the exponential rate of growth of the iterates of k . This is variously known as the degree complexity, the dynamical degree, or the algebraic entropy of k . When $b \neq 0$ and a is generic, $\delta(k_{a,b})$ is the largest root of the polynomial $x^3 - x^2 - 2x - 1$. When $b = 0$ and a is generic, $\delta(k_{a,0})$ is the largest root of $x^2 - x - 1$. The form of a map can change radically under birational equivalence: a simpler form for $k_{a,0}$ which was obtained in [BHM] made it more accessible to detailed analysis (see [BD1,3]).

A basic property is that k is reversible in the sense that $k = j \circ \iota$ is a composition of two involutions. Here we give (a birationally equivalent version of) k as a composition of involutions in a new way. This shows how $k_{a,b}$ fits naturally into a larger family of maps. Namely, for any polynomial F , we define the involutions

$$j_F(x, y) = (-x + F(y), y), \quad \iota(x, y) = \left(1 - x - \frac{x-1}{y}, -y - 1 - \frac{y}{x-1} \right),$$

and the family of birational maps is given by $k_F = j_F \circ \iota$. When F is constant, the family k_F is birationally equivalent to $k_{a,0}$, and when F is linear, k_F is equivalent to $k_{a,b}$. In this paper we determine the structure and degree complexity for the maps k_F :

Theorem 1. *Let n denote the degree of F . If n is even, then for generic parameters $\delta(k_F)$ is the largest root of the polynomial $x^2 - (n+1)x - 1$. If n is odd, then for generic parameters $\delta(k_F)$ is the largest root of $x^3 - nx^2 - (n+1)x - 1$.*

When k_F is not generic, the growth rate $\delta(k_F)$ decreases (i.e. $F \mapsto \delta(k_F)$ is lower semicontinuous in the Zariski topology). One of the interesting things about the family is to know which parameters are not generic as well as the corresponding values of $\delta(k_F)$ is decreased. The exceptional values of a for the family $k_{a,0}$, as well as the corresponding values of $\delta(k_{a,0})$, were found by Diller and Favre [DF]. Similarly, the exceptional values of (a, b) are given in [BD2]. Here we look at the maximally exceptional parameters for the case where F is cubic. These are the cubic maps with the slowest degree growth and give a 2 complex parameter family of maps which are (equivalent to) automorphisms:

* Research supported in part by the NSF

† Research supported in part by a Russian Academy of Sciences/CNRS program.

Theorem 2. *If $F(y) = ay^3 + ay^2 + by + 2$, $a \neq 0$, then k_F is an automorphism of a compact, complex surface \mathcal{Z} . Further, the degrees of k_F^n grow quadratically, and k_F is integrable.*

We will analyze the family k_F by inspecting the blowing-up and blowing-down behavior. That is, there are exceptional curves, which are mapped to points; and there are points of indeterminacy, which are blown up to curves. As was noted by Fornæss and Sibony [FS], if there is an exceptional curve whose orbit lands on a point of indeterminacy, then the degree is not multiplicative: $(\deg(k_F))^n \neq \deg(k_F^n)$. The approach we use here is to replace the original domain \mathbf{P}^2 by a new manifold \mathcal{X} . That is, we find a birational map $\varphi : \mathcal{X} \rightarrow \mathbf{P}^2$, and we consider the new birational map $\tilde{k} = \varphi \circ k_F \circ \varphi^{-1}$. There is a well defined map $\tilde{k}^* : Pic(\mathcal{X}) \rightarrow Pic(\mathcal{X})$, and the point is to choose \mathcal{X} so that the induced map \tilde{k} satisfies $(\tilde{k}^*)^n = (\tilde{k}^n)^*$. By the birational invariance of δ (see [BV] and [DF]) we conclude that $\delta(k_F)$ is the spectral radius of \tilde{k}^* . This method has also been used by Takenawa [T1–3]. The general existence of such a map \tilde{k} when $\delta(k) > 1$ was shown in [DF]. We comment that the construction of \mathcal{X} and \tilde{k} can yield further information about the dynamics of k (see, for instance, [BK] and [BD2]).

§1. The maps. Let us set $F(z) = \sum_{j=0}^n a_j z^j$ with $a_n \neq 0$. The map $k = j_F \circ \iota$ is the composition of the two involutions defined above. The map $k = [k_0 : k_1 : k_2]$ is given in homogeneous coordinates as

$$\begin{aligned} k_0 &= (x_0 x_1 - x_0^2)^n x_2 \\ k_1 &= x_0^{n-1} (x_1 - x_0)^{n+1} (x_2 + x_0) + x_2 \sum_{j=0}^n a_j (x_0 x_1 - x_0^2)^{n-j} (x_2^2 - x_0 x_1 - x_1 x_2)^j \\ k_2 &= x_2 (x_0 x_1 - x_0^2)^{n-1} (x_2^2 - x_0 x_1 - x_1 x_2). \end{aligned} \quad (1.1)$$

Each coordinate function has degree $2n+1$, which means that $\deg(k) = 2 \cdot \deg(F) + 1$. Since the jacobian of this map is $x_0^{3n-3} (x_0 - x_1)^{3n-1} x_2^2 (x_0^2 - x_0 x_1 - x_1 x_2)$ we have four exceptional curves :

$$C_1 := \{x_0 = 0\}, \quad C_2 := \{x_0 = x_1\}, \quad C_3 := \{x_2 = 0\}, \quad C_4 := \{-x_0^2 + x_0 x_1 + x_1 x_2 = 0\}.$$

When $a_0 \neq 2$, the exceptional hypersurfaces are mapped as:

$$k : C_4 \mapsto [1 : -1 + a_0 : 0] \in C_3 \quad \text{and} \quad C_1 \cup C_2 \cup C_3 \mapsto e_1. \quad (1.2)$$

The points of indeterminacy for k are

$$e_1 := [0 : 1 : 0], \quad e_2 := [0 : 0 : 1], \quad \text{and} \quad e_{01} := [1 : 1 : 0].$$

Figure 1.1 shows the relative position of the points of indeterminacy (dots with circles around them), exceptional curves, and the critical images (big dots). The information that $C_1, C_2, C_3 \rightarrow e_1$ is not drawn for lack of space.

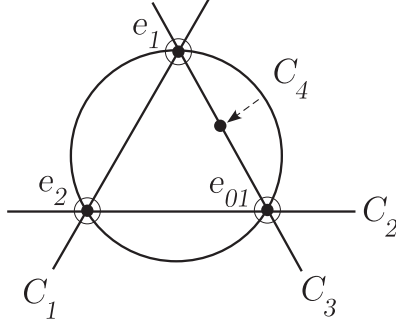


Figure 1.1. Exceptional curves and points of indeterminacy.

The sort of singularity that will be the most difficult to deal with arises from the exceptional curve $C_1 \mapsto e_1 \in C_1$. In local coordinates near e_1 , this looks like

$$k[t : 1 : y] = \left[\frac{t^n + \dots}{a_n(-y)^n + \dots} : 1 : \frac{t^{n-1} + \dots}{a_n(-y)^{n-1} + \dots} \right]. \quad (1.3)$$

For this, we will perform the iterated blowups described in §2.

The inverse map $k^{-1} = [k_0^{-1} : k_1^{-1} : k_2^{-1}]$ is given as

$$\begin{aligned} k_0^{-1} &= x_0^n x_2 (\check{F} - x_0^{n-1}(x_0 + x_1)) \\ k_1^{-1} &= (x_0 + x_2) \left(\sum_{j=0}^n a_j x_0^{n-j} x_2^j - x_0^{n-1}(x_0 + x_1) \right)^2 \\ k_2^{-1} &= x_0^{n-1} x_2 (x_0^{n-1}(x_0^2 + x_0 x_1 + x_1 x_2) - (x_0 + x_2) \check{F}) \end{aligned}$$

where $\check{F} = x_0^n F(x_2/x_0) = \sum_{j=0}^n a_j x_0^{n-j} x_2^j$. The jacobian for the inverse map is

$$x_0^{3n-3} x_2^2 (x_0^n + x_0^{n-1} x_1 - \check{F})^2 (x_0^{n+1} - (x_0 + x_2)(x_0^n + x_0^{n-1} x_1 - \check{F}))$$

The exceptional curves for k^{-1} are C'_j , $1 \leq j \leq 4$, where

$$\begin{aligned} C'_1 &= C_1, \quad C'_2 := \{x_0^n + x_0^{n-1} x_1 - \check{F} = 0\}, \quad C'_3 = C_3, \\ C'_4 &:= \{x_0^{n+1} - (x_0 + x_2)(x_0^n + x_0^{n-1} x_1 - \check{F}) = 0\}. \end{aligned}$$

$$k^{-1} : C'_1 \cup C'_3 \mapsto e_1, \quad C'_2 \mapsto e_2, \quad \text{and} \quad C'_4 \mapsto e_{01} \in C'_3, \quad (1.4)$$

§2. Blowups and local coordinate systems. In this section we discuss iterated blowups, and we explain the choices of local coordinates which will be useful in the sequel. Let $\pi : X \rightarrow \mathbf{C}^2$ denote the complex manifold obtained by blowing up the origin $e = (0, 0)$; the space is given by

$$X = \{((t, y), [\xi : \eta]) \in \mathbf{C}^2 \times \mathbf{P}^1 ; t\eta = y\xi\},$$

and π is projection to \mathbf{C}^2 . Let $E := \pi^{-1}(e)$ denote the exceptional fiber over the origin, and note that π^{-1} is well defined over $\mathbf{C}^2 - e$. The closure in X of the y -axis ($\pi^{-1}(\{t = 0\} - e)$) corresponds to the hypersurface $\{\xi = 0\} \subset X$. On the complement $\{\xi \neq 0\}$ set $u = t$ and

$\eta = y/t$. Then (u, η) defines a coordinate system on $X \setminus \{t = 0\}$, with a point being given by $((t, y), [1 : y/t]) = ((u, u\eta), [1 : \eta])$. We will use the notation $(u, \eta)_L$. On the set $t \neq 0$, the coordinate projection π is given in these coordinates as

$$\pi_L(u, \eta)_L = (u, u\eta) = (t, y) \in \mathbf{C}^2. \quad (2.1)$$

Figure 2.1 illustrates this blowup with emphasis on the relation between the point e and the lines $t = 0$ and $y = 0$ which contain it. The space X is drawn twice to show two choices of coordinate system; the dashed lines show where each coordinate system fails to be defined. The left hand copy of X shows the u, η -coordinate system in the complement of $t = 0$. The right hand side shows a different choice of coordinate; we would choose this coordinate system to work in a neighborhood of the point $p_1 := E \cap \{t = 0\}$.

In the u, η coordinate system (on the upper left side of Figure 2.1), the η -axis ($u = 0$) represents the exceptional fiber $E \cong \mathbf{P}^1$. The line $\gamma_\eta = \{(s, \eta)_L : s \in \mathbf{C}\}$ projects to the line $\{y = \eta t\} \subset \mathbf{C}^2$, and $(0, \eta)_L = E \cap \gamma_\eta$. It follows that $E \cap \{y = 0\} = (0, 0)_L$ in this coordinate system.

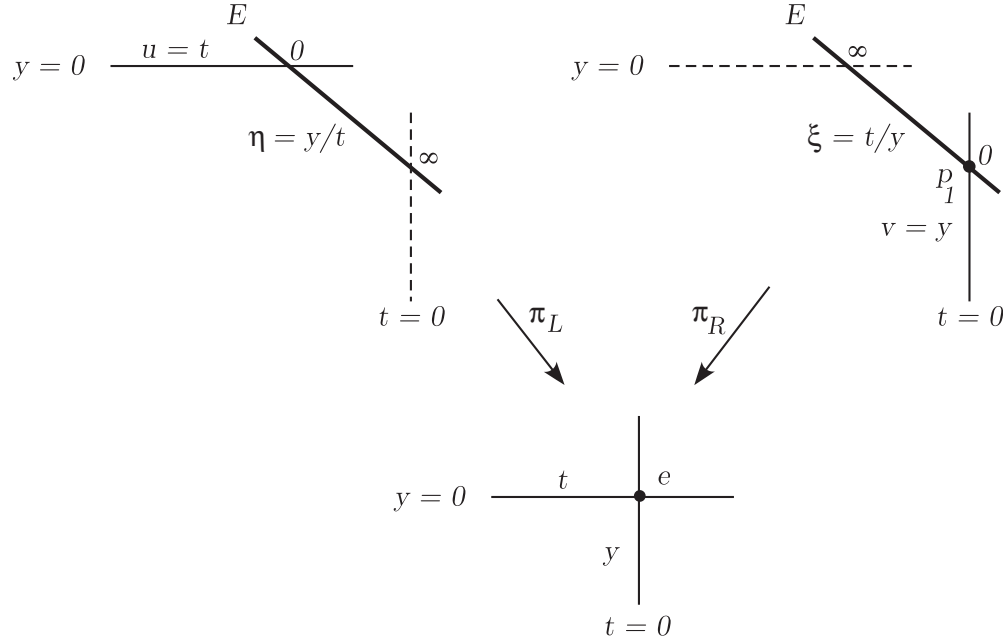


Figure 2.1. Two choices of local coordinate systems.

On the upper right side of Figure 2.1, we define a (ξ, v) -coordinate system on the complement of t -axis ($y = 0$) :

$$\pi_R : (\xi, v)_R = (t/y, y) \rightarrow (v\xi, v) \in \mathbf{C}^2. \quad (2.2)$$

The exceptional fiber E is given by ξ -axis ($v = 0$). Next we blow up $p_1 = E \cap \{t = 0\} = \{\xi = v = 0\} = (0, 0)_R$. Let P_1 denote the exceptional fiber over p_1 . The choice of a local coordinate system depends on the center of next blowup. Suppose the third blowup center is an intersection of two exceptional fibers $p_2 := E \cap P_1$. For this we are led to the (u, η) -coordinate system, as on the left side of Figure 2.1. Thus we have a local coordinate system on the complement of $\{t = 0\} \cup \{y = 0\}$;

$$(u_1, \eta_1)_1 = (t/y, y^2/t) \rightarrow (u_1, u_1\eta_1)_R \rightarrow (u_1^2\eta_1, u_1\eta_1) \in \mathbf{C}^2. \quad (2.3)$$

This (u_1, η_1) -coordinate system is defined only off the axes $(t = 0) \cup (y = 0)$; the new exceptional fiber P_1 is given by the η_1 -axis.

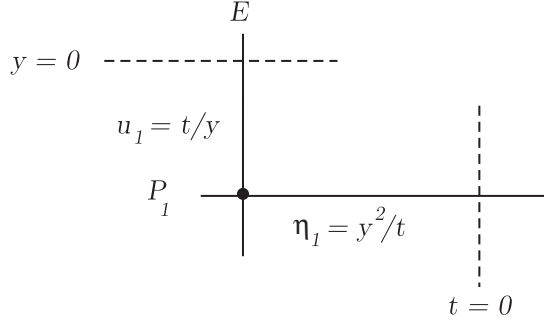


Figure 2.2. Blowup of p_1 in (u_1, η_1) -coordinates.

Now we define a sequence of iterated blowups which will let us deal with the singularity (1.3). We start with the blowup space X as in Figure 2.2, and we continue inductively for $2 \leq j \leq n$ by setting $p_j := E \cap P_{j-1}$ and letting P_j be the exceptional fiber. For each $2 \leq j \leq n$, we use the left-hand coordinate system of Figure 2.1, which corresponds to (2.1). Thus we have the coordinate projection $\pi_j : P_j \rightarrow \mathbf{C}^2$:

$$\pi_j : (u, \eta)_j \rightarrow (u^{j+1}\eta, u^j\eta) = (t, y) \in \mathbf{C}^2, \quad \pi_j^{-1}(t, y) = (u, \eta) = (t/y, y^{j+1}/t^j). \quad (2.4)$$

This coordinate system is defined off of $\{y = 0\} \cup \{t = 0\} \cup P_1 \cup \dots \cup P_{j-1}$. A point $(0, \eta = c)_j \in P_j$ is the landing point of the curve $u \mapsto (u, c)_j$ as $u \rightarrow 0$, which projects to the curve $u \mapsto (t(u) = u^{j+1}c, y(u) = u^j c) \in \mathbf{C}^2$. In Figure 2.3, the exceptional fibers P_j , $1 \leq j \leq n$ are drawn with their fiber coordinates y^{j+1}/t^j .

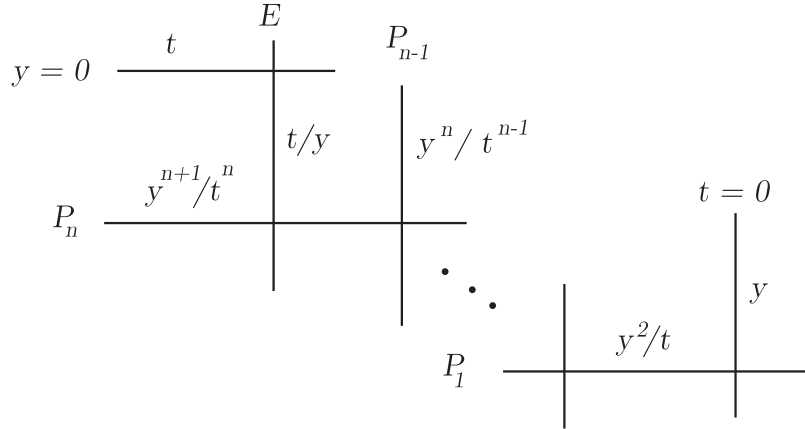


Figure 2.3. n -th iterated blowup

§3. Mappings with $n = \text{even}$. We define a complex manifold $\pi_{\mathcal{X}} : \mathcal{X} \rightarrow \mathbf{P}^2$ by blowing up points $e_1, q, p_1, \dots, p_{n-1}$ in the following order:

- (i) blow up $e_1 = [0 : 1 : 0]$ and let E_1 denote the exceptional fiber over e_1 ,
- (ii) blow up $q := E_1 \cap C_4$ and let Q denote the exceptional fiber over q ,
- (iii) blow up $p_1 := E_1 \cap C_1$ and let P_1 denote the exceptional fiber over p_1 ,

(iv) blow up $p_j := E_1 \cap P_{j-1}$ with exceptional fiber P_j for $2 \leq j \leq n-1$.

The iterated blow-up of p_1, \dots, p_{n-1} is exactly the process described in §2, so we will use the local coordinate systems defined there. That is, in a neighborhood of Q we use a $(\xi_1, v_1) = (t^2/y, y/t)$ coordinate system. For E_1 and $P_j, 1 \leq j \leq n-1$ we use local coordinate systems defined in (2.2–4). We use homogeneous coordinates by identifying a point $(t, y) \in \mathbf{C}^2$ with $[t : 1 : y] \in \mathbf{P}^2$. Let $k_{\mathcal{X}} : \mathcal{X} \rightarrow \mathcal{X}$ denote the induced map on the complex manifold \mathcal{X} . In the next few lemmas, we will show that $k_{\mathcal{X}}$ maps the exceptional fibers as shown in Figure 3.1.

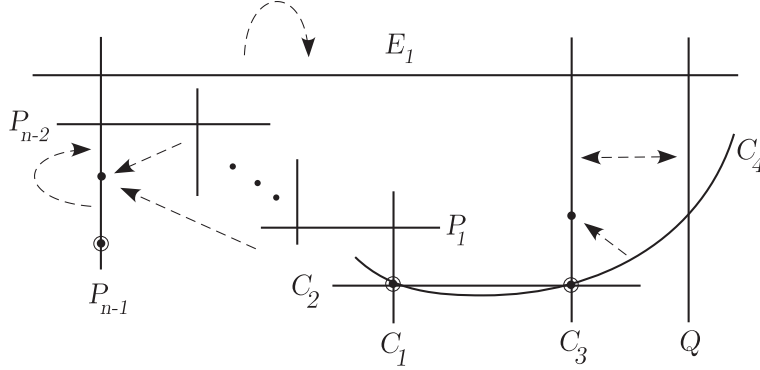


Figure 3.1. The space \mathcal{X} and the action of $k_{\mathcal{X}}$

Lemma 3.1. *Under the induced map $k_{\mathcal{X}}$, the blowup fibers E_1 and P_{n-1} are mapped to themselves :*

$$\begin{aligned} k_{\mathcal{X}} : E_1 \ni \xi &\mapsto -\xi/(\xi + 1) \in E_1 \\ P_{n-1} \ni \eta_{n-1} &\mapsto \eta_{n-1}/(1 + a_n \eta_{n-1}) \in P_{n-1}. \end{aligned} \quad (3.1)$$

Proof. First let us work on E_1 . We use the local coordinate system defined in (2.2), so a point in the exceptional fiber E_1 is $(\xi, 0)_R$. To see the forward image of E_1 we consider a nearby point $(\xi, v)_R \rightarrow (v\xi, v)$ with small v and we have $k_{\mathcal{X}}(\xi, 0)_R = \lim_{v \rightarrow 0} k_{\mathcal{X}}(\xi, v)_R$. By (1.1) we see that

$$k[v\xi : 1 : v] = [v\xi + \dots : 1 + \dots : -v(\xi + 1) + \dots]$$

where we use \dots to indicate the higher order terms in v . As in Figure 2.1, the coordinate of the landing point in E_1 is given by the ratio of t - and y -coordinates. Thus we have

$$k_{\mathcal{X}}|_{E_1} : \xi \mapsto \lim_{v \rightarrow 0} k_0/k_2 = \lim_{v \rightarrow 0} (v\xi + \dots)/(-v(\xi + 1) + \dots) = -\xi/(\xi + 1).$$

Now we determine the behavior of $k_{\mathcal{X}}$ on P_{n-1} . A fiber point $(0, \eta_{n-1}) \in P_{n-1}$ is the landing point of the arc $u \mapsto (u, \eta_{n-1})$ as $u \rightarrow 0$. To show that $k_{\mathcal{X}}$ maps P_{n-1} to P_{n-1} , we need to evaluate:

$$\lim_{u \rightarrow 0} k_{\mathcal{X}}(u, \eta_{n-1}) = \lim_{u \rightarrow 0} \pi_{n-1}^{-1} \circ k \circ \pi_{n-1}(u, \eta_{n-1}).$$

Using the formulas for π_{n-1} and π_{n-1}^{-1} in (2.4), we obtained the desired limit. \square

Now we may use similar calculations to show that $k_{\mathcal{X}} : P_j \rightarrow P_{n-1}$; we fix a point $(0, \eta_j) \in P_j$ and show the existence of the limit

$$\lim_{u \rightarrow 0} k_{\mathcal{X}}(u, \eta_j) = \lim_{u \rightarrow 0} \pi_{n-1}^{-1} \circ k \circ \pi_j(u, \eta_j).$$

Doing this, we find that the line C_1 and all blowup fibers $P_j, j = 1, \dots, n-2$ are all exceptional for both $k_{\mathcal{X}}$ and $k_{\mathcal{X}}^{-1}$. And C_2 is exceptional for $k_{\mathcal{X}}$:

$$\begin{aligned} k_{\mathcal{X}} & : C_1, C_2, P_1, \dots, P_{n-2} \mapsto 1/a_n \in P_{n-1} \\ k_{\mathcal{X}}^{-1} & : C_1, P_1, \dots, P_{n-2} \mapsto (-1)^{n-1}/a_n \in P_{n-1} \end{aligned} \tag{3.2}$$

Combining (3.1–2) it is clear that the indeterminacy locus of $k_{\mathcal{X}}$ consists of three points

$$e_2, e_{01}, \quad \text{and} \quad (-1)^{n-1}/a_n \in P_{n-1}.$$

Lemma 3.2. *If n is even, then the orbits of the exceptional curves $C_1, C_2, P_1, \dots, P_{n-2}$ are disjoint from the indeterminacy locus.*

Proof. By Lemma 3.1, the orbit of $1/a_n$ in P_{n-1} is $\{1/a_n, 1/(2a_n), 1/(3a_n), \dots\} \subset P_{n-1}$. This is disjoint from the indeterminacy locus since it does not contain point $-1/a_n$ in P_{n-1} . \square

A computation as in the proof of Lemma 3.1 shows that $k_{\mathcal{X}}$ maps $Q \leftrightarrow C_3$ according to:

$$\begin{aligned} k_{\mathcal{X}} & : Q \ni \xi_1 \mapsto [1 : a_0 - \xi_1 : 0] \in C_3, \\ C_3 & \ni [x_0 : x_1 : 0] \mapsto -x_1/x_0 \in Q. \end{aligned} \tag{3.3}$$

Lemma 3.3. *If $a_0 \neq 2/m$ for all $m > 0$ then the indeterminacy locus of $k_{\mathcal{X}}$ and the forward orbit of C_4 under the induced map $k_{\mathcal{X}}$ are disjoint. If $a_0 = 2/m$ for some $m > 0$, we have $k_{\mathcal{X}}^{2m-1}C_4 = e_{01}$.*

Proof. Since the forward image of C_4 is $[1 : -1 + a_0 : 0] \in C_3$, using (3.3) we have that $k_{\mathcal{X}}^{2m-1}C_4 = [1 : ma_0 - 1 : 0] \in C_3$. Since the unique point of indeterminacy in C_3 is e_{01} , for C_4 to be mapped to a point of indeterminacy, a_0 must satisfy $ma_0 - 1 = 1$ for some $m \geq 0$. \square

The following theorem comes directly from previous Lemmas.

Theorem 3.4. *Suppose that n is even and $a_0 \neq 2/m$ for all integers $m \geq 0$. Then no orbit of an exceptional curve contains a point of indeterminacy.*

Let us recall the Picard group $Pic(\mathcal{X})$, which is the set of all divisors in \mathcal{X} , modulo linear equivalence, which means that $D_1 \sim D_2$ if $D_1 - D_2$ is the divisor of a rational function. $Pic(\mathbf{P}^2)$ is 1-dimensional and generated by the class of any line (hyperplane) H , and a basis of $Pic(\mathcal{X})$ is given by the class of a general hyperplane $H_{\mathcal{X}} := \pi^*H$, together with all of the blowup fibers $E_1, Q, P_1, \dots, P_{n-1}$. If r is a rational function on \mathcal{X} , then the pullback $k_{\mathcal{X}}^*r := r \circ k_{\mathcal{X}}$ is just the composition. To pull back a divisor, we just pull back its defining functions. This gives the pullback map $k_{\mathcal{X}}^* : Pic(\mathcal{X}) \rightarrow Pic(\mathcal{X})$. Thus from (3.1-2) we see

that the pullback of E_1 is E_1 and the pulling back of most of basis elements are trivial, that is $k_{\mathcal{X}}^*P_j = 0$ for all $j = 1, \dots, n-2$.

Next we pull back $H_{\mathcal{X}}$. Since k has degree $2n+1$ we have $k^*H = (2n+1)H$ in $Pic(\mathbf{P}^2)$. Now we pull back by $\pi_{\mathcal{X}}^*$ to obtain:

$$(2n+1)H_{\mathcal{X}} = \pi_{\mathcal{X}}^*(2n+1)H = \pi_{\mathcal{X}}^*(k^*H). \quad (3.4)$$

A line is given by $\{h := \alpha_0x_0 + \alpha_1x_1 + \alpha_2x_2 = 0\}$, so k^*H is the divisor defined by $h \circ k = \sum_j \alpha_j k_j$. To write this divisor as a linear combination of basis elements $H_{\mathcal{X}}, E_1, Q, P_1, \dots, P_{n-1}$, we need to check the order of vanishing of $h \circ k$ at all of these sets. Let us start with the coordinate system $\pi_{\mathcal{X}}(\xi, v) = [v\xi : 1 : v]$ near E_1 , defined in §2. Using the expression for k given in §1 we see that $\alpha_0k_0 + \alpha_1k_1 + \alpha_2k_2$ vanishes to order n in v . It follows that $\pi_{\mathcal{X}}^*k^*H$ vanishes at E_1 with multiplicity n . Similar computations for all other basis elements gives us $\pi_{\mathcal{X}}^*k^*H = k_{\mathcal{X}}^*H_{\mathcal{X}} + nE_1 + (n+1)Q + (n+1)\sum_j jP_j$. Combining with (3.4) we have

$$k_{\mathcal{X}}^*H_{\mathcal{X}} = (2n+1)H_{\mathcal{X}} - nE_1 - (n+1)Q - (n+1)\sum_{j=1}^{n-1} jP_j. \quad (3.5)$$

Similarly, we obtain:

$$\begin{aligned} k_{\mathcal{X}}^* : Q &\mapsto H_{\mathcal{X}} - E_1 - Q - P_1 - 2P_2 - \dots - (n-1)P_{n-1} \\ P_{n-1} &\mapsto 2H_{\mathcal{X}} - E_1 - Q - P_1 - 2P_2 - \dots - (n-1)P_{n-1}. \end{aligned} \quad (3.6)$$

Theorem 1: $n = \text{even}$. Suppose $F(z) = \sum_{j=1}^n a_j z^j$ is an even degree polynomial associated with J_F . If $a_0 \neq 2/m$ for any positive integer m , then the degree complexity is the largest root of the quadratic polynomial $x^2 - (n+1)x - 1$.

Proof. Since P_1, \dots, P_{n-2} are mapped to 0 under the action on cohomology, it suffices to consider the action restricted to $H_{\mathcal{X}}, E_1, Q$, and P_{n-1} . By (3.5,6) the matrix representation of $k_{\mathcal{X}}^*$, restricted to the ordered basis $\{H_{\mathcal{X}}, E_1, Q, P_{n-1}\}$, is

$$\begin{pmatrix} 2n+1 & 0 & 1 & 2 \\ -n & 1 & -1 & -1 \\ -n-1 & 0 & -1 & -1 \\ -n^2+1 & 0 & -n+1 & -n+1 \end{pmatrix}.$$

The characteristic polynomial is $x(x-1)(x^2 - (n+1)x - 1)$. □

§4. Mappings with $n = \text{odd}$. Let us start with the space \mathcal{X} from §3. When n is odd, we see from (3.2) that the image of all exceptional lines of $k_{\mathcal{X}}$ coincide with a point of indeterminacy in $p_n \in P_{n-1}$. Let $\pi_{\mathcal{Y}} : \mathcal{Y} \rightarrow \mathbf{P}^2$ be the complex manifold obtained by blowing up \mathcal{X} at the point p_n , and let P_n denote the exceptional fiber over p_n . In the u_{n-1}, η_{n-1} coordinate system, p_n has coordinate $(0, 1/a_n)_{n-1}$. Thus, at P_n , we use the coordinate projection:

$$\pi_n : \mathcal{Y} \ni (u, \eta)_n \rightarrow (u^n(u\eta + 1/a_n), u^{n-1}(u\eta + 1/a_n)) \in \mathbf{C}^2.$$

Most computations in the previous section remain valid for n odd. Thus Lemma 3.3, (3.1) and (3.3) are still valid for the induced map $k_{\mathcal{Y}} : \mathcal{Y} \rightarrow \mathcal{Y}$. Under $k_{\mathcal{Y}}$ curves $C_1, C_2, P_1, \dots, P_{n-3}$ are still exceptional:

$$\begin{aligned} k_{\mathcal{Y}} & : C_1, C_2, P_1, \dots, P_{n-3} \mapsto -a_{n-1}/a_n^2 \in P_n \\ k_{\mathcal{Y}}^{-1} & : C_1, P_1, \dots, P_{n-3} \mapsto (a_{n-1} - (n-1)a_n)/a_n^2 \in P_n. \end{aligned} \quad (4.1)$$

The blowup fibers P_n and P_{n-2} form a two cycle, $k_{\mathcal{Y}} : P_n \leftrightarrow P_{n-2}$ and P_{n-1} is mapped to itself as before. It follows that the points of indeterminacy for $k_{\mathcal{Y}}$ are e_2, e_{01} and $(a_{n-1} - (n-1)a_n)/a_n^2 \in P_n$. For all $m \geq 0$, we have

$$k_{\mathcal{Y}}^{2m} : P_n \ni -a_{n-1}/a_n^2 \mapsto (2m(n-1)a_n - (4m+1)a_{n-1})/a_n^2 \in P_n \quad (4.2)$$

As a consequence of (4.1) and (4.2) we have:

Lemma 4.1. *If n is odd, and if*

$$2a_{n-1} \neq (n-1)a_n, \quad (4.3)$$

then the forward orbits of $C_1, C_2, P_1, \dots, P_{n-3}$ under $k_{\mathcal{Y}}$ do not contain any point of indeterminacy.

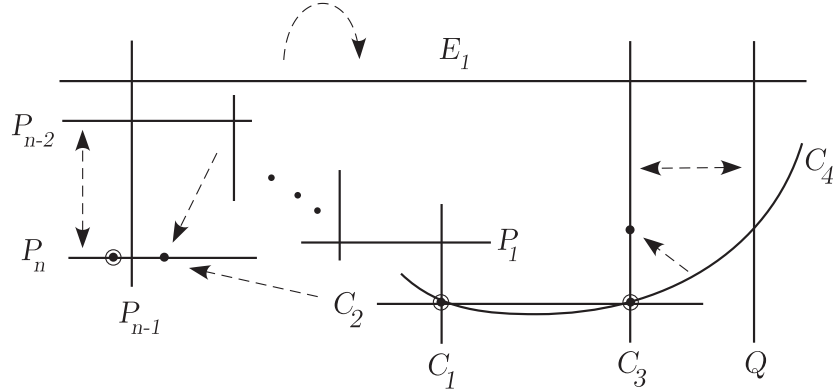


Figure 4.1. The space \mathcal{Y} and the action of $k_{\mathcal{Y}}$.

Combining Lemma 3.3 and 4.1 we have

Theorem 4.2. *Suppose that n is odd, $a_0 \neq 2/m$ for all $m > 0$, and $a_{n-1} \neq (n-1)a_n/2$. Then the forward orbits of exceptional curves do not contain any points of indeterminacy.*

To determine $k_{\mathcal{Y}}$, we use the basis $\{H_{\mathcal{Y}}, E_1, Q, P_1, \dots, P_n\}$ for $Pic(\mathcal{Y})$. Now the exceptional lines $C_1, C_2, P_1, \dots, P_{n-2}$ are mapped to P_n . Let $\{C_1\} \in Pic(\mathcal{Y})$ denote the class of the strict transform of C_1 , i.e., the closure in \mathcal{Y} of $\pi_{\mathcal{Y}}^{-1}(C_1 - \text{centers of blowup})$. (The curve C_2 does not pass through any center of blowup, so with the same notation we have $\{C_2\} = H_{\mathcal{Y}} \in Pic(\mathcal{Y})$.) In order to write $\{C_1 = (x_0 = 0)\}$ in terms of our basis, we note first that $\pi_{\mathcal{Y}}^{-1}C_1 = C_4 \cup E_1 \cup Q \cup P_1 \cup \dots \cup P_{n-1}$, i.e., the pullback function $x_0 \circ \pi_{\mathcal{Y}}$ vanishes on all of these curves. Thus we have to compute the multiplicities of vanishing. At P_{n-1} , for instance, we consider the (u_{n-1}, η_{n-1}) coordinate system defined in (2.4), and we see that

$k_{\mathcal{Y}}^*x_0$ vanishes to order n at $P_{n-1} = (u_{n-1} = 0)$. Similarly we can compute the multiplicities for $E_1, Q, P_1, \dots, P_{n-2}$ and P_n , so

$$H_{\mathcal{Y}} = \pi_{\mathcal{Y}}^*C_1 = \{C_1\} + E_1 + Q + 2P_1 + 3P_2 + \dots + nP_{n-1} + nP_n.$$

It follows that

$$k_{\mathcal{Y}}^*P_n = \{C_1\} + \{C_2\} + \sum_{j=1}^{n-2} P_j = 2H_{\mathcal{Y}} - E_1 - Q - \sum_{j=1}^{n-2} jP_j - nP_{n-1} - nP_n.$$

For the rest of basis entries we have

$$\begin{aligned} k_{\mathcal{Y}}^* : H_{\mathcal{Y}} &\mapsto (2n+1)H_{\mathcal{Y}} - nE_1 - (n+1)Q - (n+1)\sum_{j=1}^{n-1} jP_j - n^2P_n \\ Q &\mapsto H_{\mathcal{Y}} - E_1 - Q - P_1 - 2P_2 - \dots - (n-1)P_{n-1} - (n-1)P_n, \\ E_1 &\mapsto E_1, \quad P_{n-2} \mapsto P_n, \quad \text{and} \quad P_{n-1} \mapsto P_{n-1}. \end{aligned}$$

Theorem 1: $n = \text{odd}$. *If $a_0 \neq 2/m$ for all $m > 0$, then the degree complexity is the largest root of the cubic polynomial $x^3 - nx^2 - (n+1)x - 1$.*

Proof. The classes of the exceptional fibers P_1, \dots, P_{n-3} are all mapped to 0, and exceptional fibers E_1 and P_{n-1} are simply interchanged. It follows that to get the spectral radius of $k_{\mathcal{Y}}^*$ we only need to consider 4×4 matrix with ordered basis $\{H_{\mathcal{X}}, Q, P_{n-2}, P_n\}$ and the spectral radius is given by the largest root of $x^3 - nx^2 - (n+1)x - 1$. \square

§5. Degree 3: a family of automorphisms. Let us consider the 2 parameter family of maps $k = j_F \circ \iota$ where $F(z) = az^3 + az^2 + bz + 2$ with $a \neq 0$. We consider the complex manifold $\pi_{\mathcal{Z}} : \mathcal{Z} \rightarrow \mathbf{P}^2$ obtained by blowing up 6 points $e_2, e_{01}, p_4, p_5, p_6, r$ in the complex manifold \mathcal{Y} constructed in §4. As we construct the blowups, we will let $E_2, E_{01}, P_4, P_5, P_6$ and R denote the exceptional fibers over $e_2, e_{01}, p_4, p_5, p_6$, and r respectively. Specifically, we blow up e_2 and e_{01} and then:

$$\begin{aligned} p_4 &:= -1/a \in P_3, \quad p_5 := (2-b)/a \in P_4, \\ p_6 &:= (2b-2-a)/a^2 \in P_5, \quad \text{and} \quad r := 0 \in E_2 \cap \{x_1 = 0\}. \end{aligned}$$

We define the local coordinate system in a similar way we define local coordinates in §2. Using these local coordinates we can easily verify that under the induced map $k_{\mathcal{Z}}$ we have

$$C_1 \rightarrow P_4 \rightarrow C_1, \quad E_2 \rightarrow P_5 \rightarrow E_2, \quad C_4 \rightarrow E_{01} \rightarrow C'_4, \quad \text{and} \quad C_2 \rightarrow P_6 \rightarrow R \rightarrow C'_2$$

and all mappings are dominant and holomorphic.

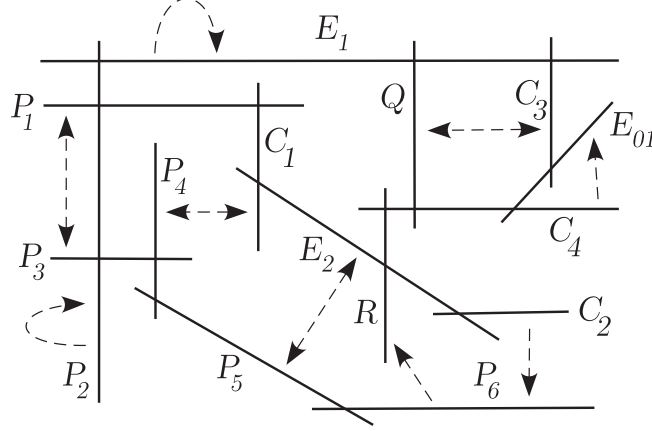


Figure 5.1. The space \mathcal{Z} and action of $f_{\mathcal{Z}}$.

For example, let us consider E_2 . We may use coordinates w, ζ which are mapped by $\pi_{E_2} : (x, \zeta) \rightarrow [w : w\zeta : 1] \in \mathbf{P}^2$. Thus $E_2 = (w = 0)$ is given by ζ -axis in this coordinate system and by considering $\lim_{w \rightarrow 0} \pi_{P_5}^{-1} \circ k \circ \pi_{E_2}(w, \zeta)$ we find:

$$k_{\mathcal{Z}} : E_2 \ni \zeta \mapsto (2b - a - \zeta - 1)/a^2 \in P_5.$$

The mapping among the exceptional fibers is shown in Figure 5.1. What is not shown is that $R \rightarrow C'_2$ and $E_{01} \rightarrow C'_4$

Theorem 5.1. *Suppose $F(z) = az^3 + az^2 + bz + 2$ with $a \neq 0$. Then the induced map $k_{\mathcal{Z}}$ is biholomorphic.*

Proof. Since $k_{\mathcal{Z}}$ and $k_{\mathcal{Z}}^{-1}$ have no exceptional hypersurface, indeterminacy locus for $k_{\mathcal{Z}}$ is empty. It follows that $k_{\mathcal{Z}}$ is an automorphism of \mathcal{Z} . \square

Repeating the argument in previous two sections, we have that $k_{\mathcal{Z}}^*$ acts on each basis element as follows :

$$\begin{aligned} H_{\mathcal{Z}} &\mapsto 7H_{\mathcal{Z}} - 3E_1 - 4P_1 - 8P_2 - 9P_3 - 10P_4 - 10P_5 - 10P_6 - 3E_2 - 6R - 4Q - 4E_{01}, \\ E_1 &\mapsto E_1, \quad P_1 \mapsto P_3 \mapsto P_1, \quad \text{and} \quad P_2 \mapsto P_2, \\ P_4 &\mapsto H_{\mathcal{Z}} - E_1 - 2P_1 - 3P_2 - 3P_3 - 3P_4 - 3P_5 - 3P_6 - E_2 - R - Q, \\ P_5 &\mapsto E_2, \quad P_6 \mapsto H_{\mathcal{Z}} - E_2 - R - E_{01}, \quad E_2 \mapsto P_5, \quad \text{and} \quad E_{01} \mapsto P_6, \\ Q &\mapsto H_{\mathcal{Z}} - E_1 - P_1 - 2P_2 - 2P_3 - 2P_4 - 2P_5 - 2P_6 - Q - E_{01}, \\ E_{01} &\mapsto 2H_{\mathcal{Z}} - E_1 - P_1 - 2P_2 - 2P_3 - 2P_4 - 2P_5 - 2P_6 - E_2 - 2R - 2Q - E_{01}. \end{aligned}$$

Theorem 5.2. *Suppose $F(z) = az^3 + az^2 + bz + 2$ with $a \neq 0$. Then the degree of $k^n = k \circ \dots \circ k$ grows quadratically, and k is integrable.*

Proof. All the eigenvalues of the characteristic polynomial of $k_{\mathcal{Z}}^*$ have modulus one. The largest Jordan block in the matrix representation of $k_{\mathcal{Z}}^*$ is a 3×3 block corresponding to the eigenvalue 1. Thus the growth rate of the powers of the matrix is quadratic.

Integrability follows from more general results: Gizatullin [G] showed that if the growth rate is quadratic, then there is an invariant fibration by elliptic curves. In this case, we can

give an explicit invariant. If we define $\phi = \phi_1/\phi_2$ to be the quotient of the following two polynomials;

$$\begin{aligned}\phi_1[x_0 : x_1 : x_2] &= x_0^2 x_2^2, \\ \phi_2[x_0 : x_1 : x_2] &= -2x_0^4 + 4x_0^3 x_1 - (2 + a)x_0^2 x_1^2 + 2ax_1 x_2^2(x_0 + x_2) - 2b(x_0^3 x_2 - x_0^2 x_1 x_2),\end{aligned}$$

then $\phi \circ k = \phi$. □

References

- [A1] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa, S. Hassani and J-M. Maillard, From Yang-Baxter equations to dynamical zeta functions for birational transformations. *Statistical physics on the eve of the 21st century*, 436–490, Ser. Adv. Statist. Mech., 14, World Sci. Publishing, River Edge, NJ, 1999.
- [A2] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa, S. Hassani and J-M. Maillard, Rational dynamical zeta functions for birational transformations. *Physica A* 264 (1999) pp. 264–293. [chao-dyn/9807014](#).
- [A3] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa, S. Hassani and J-M. Maillard, Topological entropy and complexity for discrete dynamical systems, *Phys. Lett. A* 262 (1999) pp. 44–49. [chao-dyn/9806026](#).
- [A4] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa and J-M. Maillard, Growth complexity spectrum of some discrete dynamical systems. *Phys. D* 130 (1999), no. 1-2, 27–42.
- [A5] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa and J-M. Maillard, Real topological entropy versus metric entropy for birational measure-preserving transformations. *Phys. D* 144 (2000), no. 3-4, 387–433.
- [A6] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa, S. Hassani and J-M. Maillard, Real Arnold complexity versus real topological entropy for birational transformations. *J. Phys. A* 33 (2000), no. 8, 1465–1501.
- [A7] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa, S. Hassani and J-M. Maillard, Topological entropy and Arnold complexity for two-dimensional mappings. *Phys. Lett. A* 262 (1999), no. 1, 44–49.
- [A8] N. Abarenkova, J-C. Anglès d'Auriac, S. Boukraa, and J-M. Maillard, Elliptic curves from finite order recursions or non-involutive permutations for discrete dynamical systems and lattice statistical mechanics, *The European Physical Journal B.* (1998), pp. 647–661.
- [BD1] E. Bedford and J. Diller, Real and complex dynamics of a family of birational maps of the plane: the golden mean subshift, *Amer. J. Math.* 127 (2005), no. 3, 595–646.
- [BD2] E. Bedford and J. Diller, Dynamics of a Two Parameter Family of Plane Birational Maps: Maximal entropy, *J. of Geom. Analysis*, 16 (2006), no. 3, 409–430.
- [BD3] E. Bedford and J. Diller, Real dynamics of a family of plane birational maps: trapping regions and entropy zero, [arXiv.math/0609113](#).
- [BK] E. Bedford and KH Kim, Dynamics of rational surface automorphisms: Linear fractional recurrences. [arXiv:math/0611297](#)
- [BV] M. Bellon and C. Viallet, Algebraic entropy. *Commun. Math. Phys.* 204, 425–437 (1999)

- [BHM] S. Boukraa, S. Hassani and J-M. Maillard, Product of involutions and fixed points, Alg. Rev. Nucl. Sci., Vol. 2 (1998), 1–16.
- [BM] S. Boukraa and J-M. Maillard, Factorization properties of birational mappings, Physica A **220** (1995), pp. 403–470.
- [BMR1] S. Boukraa, J-M. Maillard, and G. Rollet, Almost integrable mappings, Int. J. Mod. Phys. B8 (1994), pp. 137–174.
- [BMR2] S. Boukraa, J-M. Maillard, and G. Rollet, Integrable mappings and polynomial growth, Physica A **209**, 162-222 (1994).
- [BMR3] S. Boukraa, J-M. Maillard, and G. Rollet, Determinantal identities on integrable mappings, Int. J. Mod. Phys. **B8** (1994), pp. 2157–2201.
- [DF] J. Diller and C. Favre, Dynamics of bimeromorphic maps of surfaces, Amer. J. Math., 123 (2001), 1135–1169.
- [FS] J-E. Fornæss and N. Sibony, Complex dynamics in higher dimension, II, Modern Methods in Complex Analysis, Ann. of Math. Studies, vol. 137, Princeton U. Press, 1995, p. 135–182.
- [G] M. Gizatullin, Rational G-surfaces, Izv. Akad. Nauk SSSR Ser. Mat. (44) 1980 110–144.
- [T1] T. Takenawa, A geometric approach to singularity confinement and algebraic entropy, J. Phys. A: Math. Gen. 34 (2001) L95–L102.
- [T2] T. Takenawa, Discrete dynamical systems associated with root systems of indefinite type, Commun. Math. Phys., 224, 657–681 (2001).
- [T3] T. Takenawa, Algebraic entropy and the space of initial values for discrete dynamical systems, J. Phys. A: Math. Gen. 34 (2001) 10533–10545.

E. Bedford: bedford@indiana.edu
 T. Tuyen: ttuyen@indiana.edu
 Department of Mathematics
 Indiana University
 Bloomington, IN 47405 USA

K. Kim: kim@math.fsu.edu
 Department of Mathematics
 Florida State University
 Tallahassee, FL 32306 USA

N. Abarenkova: nina@pdmi.ras.ru
 Laboratory of Mathematical Problems of Physics
 Petersburg Department of Steklov Institute of Mathematics
 27, Fontanka, 191023, St. Petersburg, Russia

J-M. Maillard: maillard@lptmc.jussieu.fr
 Lab. de Physique Théorique et de la Matière Condensée
 Université de Paris 6, Tour 24
 4ème étage, case 121, 4, Place Jussieu
 75252 Paris Cedex 05, France